Electrons don't like each other : Impact of quantum dot shape on readout and exchange interactions

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Take home messages

-Unless you have very small and isotropic dots the single-particle picture is the wrong way of thinking about multiparticle states. Formation of Wigner molecules.

-ST splitting decreases exponentially with the dot length and the anisotropy only makes it faster

-Si dots are clearly more affected than GaAs and Ge dots

-Low-lying triplet states complicate PSB readout

-Low-lying triplet states limit exchange coupling

Intro: current status

- -High single-qubit gate fidelities have been demonstrated in both electrons and hole qubits. See
- Nature Nanotechnology **13**, 102-106 (2018)
- Nature volume **569**, pages 532–536 (2019)
- Nature Communications volume 11, Article number: 4144 (2020)
- Two-qubit gates beyond QEC threshold. See arXiv:2107.00628
- -Working four-qubit processor with holes *Nature* **591**, 580–585 (2021)

Multi-particle states. Why do they matter?

-Multi-electron states may screen part of the charge noise

Validity of the single-particle description and charge noise resilience for multielectron quantum dots Phys. Rev. B **91**, 155425 (2015) **All-Microwave Control and Dispersive Readout of Gate-Defined Quantum Dot Qubits in Circuit Quantum Electrodynamics** Phys. Rev. Lett. **122**, 206802 (2019)

-Enhanced dipolar coupling

Coherent spin control of s-, p-, d- and f-electrons in a silicon quantum dot Nature Communications **11**, Article number: 797 (2020)

-Multiple coherent oscillations **Coherent Control and Spectroscopy of a Semiconductor Quantum Dot Wigner Molecule** Phys. Rev. Lett. **127**, 127701 (2021)

Goal: understand Coulomb effects in QD and its consequences on SC qubits

-Focus on 2-particle problem (holes or electrons)

-We will make use of Hooke's model: Coulomb repulsion + harmonic confinement

-2D and single-band approximation

-General model that can be applied to different materials

Hooke's atom model

$$
H = \sum_{i=1,2} \frac{\mathbf{p}_i^2}{2m} + \frac{1}{2} m (\omega_x^2 x_i^2 + \omega_y^2 y_i^2) + \frac{e^2}{4\pi \varepsilon r}
$$

We take x as the weakest confinement direction

$$
\ell_x = \sqrt{\hbar/(m\omega_x)}
$$

$$
E_{orb} = \hbar \omega_x
$$

 $E_{ee} = e^2/(4\pi\varepsilon\ell_x)$

 $=$ $H_r = \frac{\mathbf{p}_r^2}{m} + \frac{1}{4}m(\omega_x^2 x^2 + \omega_y^2 y^2) + \frac{e^2}{4\pi\varepsilon r}$ Wigner ratio: $\lambda_W=E_{ee}/E_{orb}$ $\alpha = \omega_y/\omega_x$ Anisotropy: $(x', y') = (x/\ell_x, y/\ell_x)$

 $H_R = \frac{\mathbf{p}_R^2}{4m} + m(\omega_x^2 X^2 + \omega_y^2 Y^2)$

$$
H_r' = - \partial_{x'}^2 - \partial_{y'}^2 + \frac{1}{4} (x'^2 + \alpha^2 y'^2) + \frac{\lambda_W}{|r'|}
$$

Spectral compression & single-particle approximation breakdown

 $\Delta_{ST} < \hbar \omega_x$

ST-splitting is reduced

-The Coulomb repulsion induces a reduction of the ST-splitting

-The higher the interaction (larger dot) the larger the reduction

-The anisotropy quickly enhances this effect

Silicon is the most affected

Comparison with realistic devices

-Poisson solver for the electrostatics

-4KP band calculation for the single-hole simulations

-TB for the singleelectron simulations

-Full CI to extend to 2 particle states

Simple model agrees with simulations

Singlet-triplet splitting decays fast with gate length

-2-hole results show fast ST decay with L

-Strengthening the confinement with gates mitigates the decay

What about valleys?

-Used TB to obtain two-electron spectrum with valley states

-The valley triplet is mostly insensitive to molecularization even under disorder

-If the dot is too large, the orbital triplet will go lower than the valley triplet anyways

PSB readout: what happens if I have a low triplet state?

-Using PSB for readout requires going to $(1,1)/(0,2)$ anticrossing

-If the qubit is in S state, a charge transition occurs

-If the qubit is in the T state, there should be no charge signature

-A low triplet may yield false readouts of singlet states

Readout may be hampered even at optimal positions

 $P_{\rm opt} = P_{S(1,1)\rightarrow S(0,2)}(\varepsilon_{\rm opt}) - P_{T(1,1)\rightarrow T(0,2)}(\varepsilon_{\rm opt})$

Effects on the exchange interaction

-The exchange coupling is just $J = E_T - E_S$

-The singlet state is bent with respect to the triplet state by the presence of excited singlets

-Low-lying triplets bend the triplet energy in a similar way, limiting J

$$
J = J_{\infty} - \frac{1}{2} \left(\varepsilon - \Delta_{ST} + \sqrt{(\varepsilon - \Delta_{ST})^2 + 4\tau_T^2} \right)
$$

What about the SOP?

-At the SOP the qubits are insensitive to charge noise

Impact of dot connectivity: face to face vs lateral

-In F2F the $T(1,1)$ configuration has a negligible tunnel coupling to the T(0,2): Bending is limited

-In the lateral configuration, $T(1,1)$ loves to go to $T(0,2)$

Highest bending

Dashed for T(0,2) at infinity Solid for realistic T(0,2)

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Smaller + isotropic dots are the easiest way to avoid this